

From the microstructure to condensed meta-material models

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Sorrento

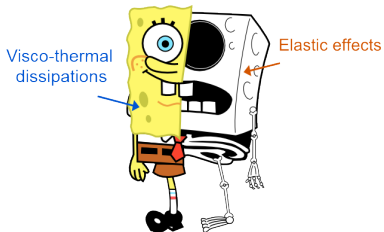
SAPEM' 23

常熟

- 1 Acoustical modeling of porous media
- 2 Micro-Macro approaches
- 3 Condensed models
- 4 Conclusion

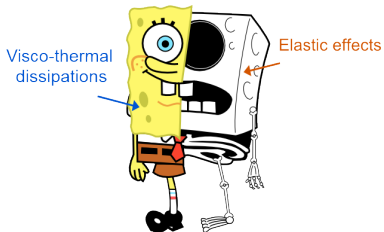
Hypothesis

- Zwikker & Kosten (1949) introduce the hypothesis of a decoupling between the visco-inertial and (micro)thermal effects.
- Biot (1956) considers visco-inertio-thermal effects & elastical effects and their coupling.

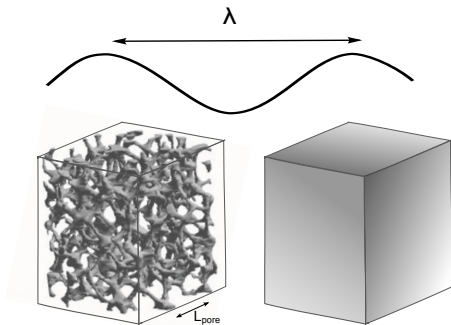


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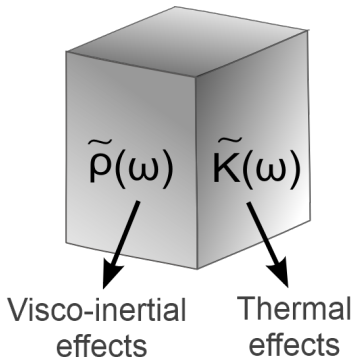
Equivalent fluid



- A porous material may be viewed as an **equivalent fluid** only if there is an order of magnitude between the observation wavelength λ and the cell characteristic size ℓ .

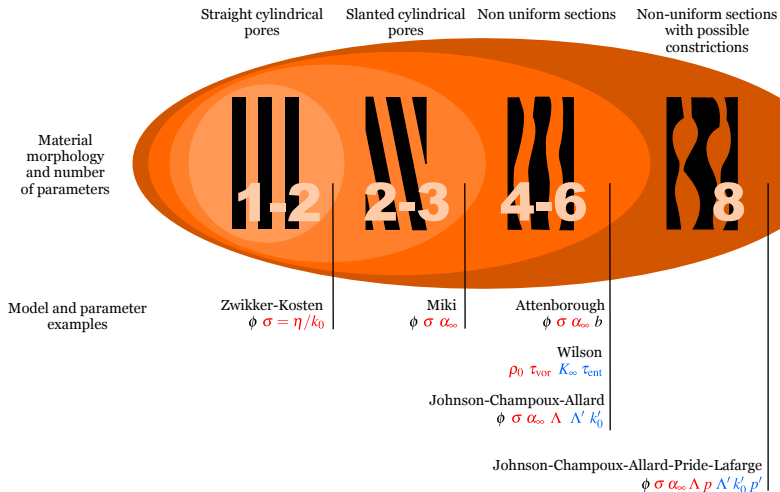
Equivalent fluid

- Visco-thermal dissipation are taken into account using the dynamic **mass density** $\tilde{\rho}(\omega)$ and the dynamic **bulk modulus** $\tilde{K}(\omega)$.



These quantities are complex and frequency dependent.

Models of visco-thermal dissipation



Macroscopic parameters

- ϕ : open porosity
- σ : static air flow resistivity
- Λ : viscous characteristic length
- Λ' : thermal characteristic length
- α_∞ : high frequency limit of tortuosity
- k'_0 : static thermal permeability

- α_0 : static viscous tortuosity
- α'_0 : static thermal tortuosity
- k_0 : static viscous permeability ($= \eta/\sigma$)

Analytical computation : perforated plate

- A **perforated plate** may be viewed as a **porous material**.
- **Macroscopic parameters:**

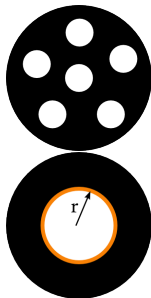
ϕ : perforation rate

$$\Lambda = \Lambda' = r$$

$$\sigma = \frac{8\eta}{\phi r^2}$$

$$\alpha_\infty = 1 + \frac{n\epsilon}{L}$$

Atalla N, Sgard F. "Modeling of perforated plates and screens using rigid frame porous models", J. Sound Vib. 303 (2007).







L : Thickness of the plate

ϵ : Length correction

n : Factor depending on the nature of upstream and downstream materials

Analytic computation: constant cross-section

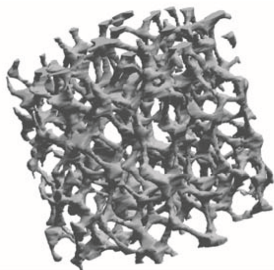
SOUND PROPAGATION THROUGH SOME MATERIALS WITH CYLINDRICAL PORES OF CONSTANT CROSS-SECTIONS	Dynamic mass density $\tilde{\rho}_{eq}$	Dynamic bulk modulus \tilde{K}_{eq}
<p>Circular (radius r)</p> 	$\frac{\rho_0}{\phi} \times 1 / \left[1 - \frac{2}{\beta\sqrt{-j}} \frac{J_1(\beta\sqrt{-j})}{J_0(\beta\sqrt{-j})} \right]$ <p>with $\beta = \sqrt{\frac{\omega \rho_0}{\eta}}$ and r the hydraulic radius equals to</p>	$\frac{P_0}{\phi} \times \gamma / \left[1 + (\gamma - 1) \frac{2}{\beta\sqrt{-j}} \frac{J_1(\beta\sqrt{-j})}{J_0(\beta\sqrt{-j})} \frac{Pr}{Pr} \right]$ <p>Formulas from C. Zwikker & C. W. Kosten 1949 after work by J. W. Strutt (lord Rayleigh).</p>
<p>Slit (width $2a$)</p> 	$\frac{\rho_0}{\phi} \times 1 / \left[1 - \frac{\tanh(\beta\sqrt{j})}{\beta\sqrt{j}} \right]$ <p>with $\beta = \sqrt{\frac{\omega \rho_0}{\eta}}$ and a the hydraulic radius equals to</p>	$\frac{P_0}{\phi} \times \gamma / \left[1 + (\gamma - 1) \frac{\tanh(\beta\sqrt{j})}{\beta\sqrt{j}} \frac{Pr}{Pr} \right]$ <p>Formulas attributed to M. Biot.</p>
<p>Equilateral triangle (length l)</p> 	$\frac{\rho_0}{\phi} \times \epsilon^2 / \left[\epsilon^2 - 3\epsilon \coth(\epsilon) + 3 \right]$ <p>with $\epsilon = \frac{\sqrt{3}}{4} \sqrt{\frac{j\omega \rho_0}{\eta}} l$</p>	$\frac{P_0}{\phi} \times \gamma / \left(\gamma - \frac{\gamma - 1}{Pr} \epsilon^2 \left[Pr \epsilon^2 - 3\sqrt{Pr} \epsilon \coth(\sqrt{Pr} \epsilon) + 3 \right] \right)$ <p>M. R. Stinson & Y Champoux. J. Acoust. Soc. Am. 91, 685-695, 1992.</p>
<p>Rectangle ($2a \times 2b$)</p> 	$\frac{\rho_0}{\phi} \times \frac{\eta \omega^2 l^2}{4j\omega \rho_0} / \left[\sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{1}{\alpha_m^2 \beta_n^2 (\alpha_m^2 + \beta_n^2 + j\omega \rho_0 / \eta)} \right]$ <p>with $\alpha_m = (m + 1/2)\pi/a$ and $\beta_n = (n + 1/2)\pi/b$</p>	$\frac{P_0}{\phi} \times \gamma / \left(\gamma - (\gamma - 1) \frac{4j\omega \rho_0 Pr}{\eta \omega^2 l^2} \left[\sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{1}{\alpha_m^2 \beta_n^2 (\alpha_m^2 + \beta_n^2 + j\omega \rho_0 / \eta)} \right] \right)$ <p>M. R. Stinson, J. Acoust. Soc. Am. 89, 550-558, 1991. H. Roh, W. P. Arnott, J. M. Sabatier & R. Raspet, J. Acoust. Soc. Am. 89, 2617-2624, 1991.</p>
<p>ρ_0 Density of air at rest kg.m^{-3} η Dynamic viscosity of air N.s.m^{-2} γ Ratio of specific heat P_0 Static/Atmospheric pressure Pa Pr Prandtl's number ω Pulsation (angular frequency) s^{-1} ϕ Open porosity of the material</p>		<p>$j = \sqrt{-1}$ J_n Bessel function of the first kind</p> <p>With a $\exp(+j\omega t)$ time convention, characteristic impedance and wavenumber are deduced from: $Z_c = \sqrt{\tilde{\rho}_{eq} \tilde{K}_{eq}}$ and $k_c = \omega \sqrt{\tilde{\rho}_{eq} / \tilde{K}_{eq}}$</p>

Micro-macro approaches

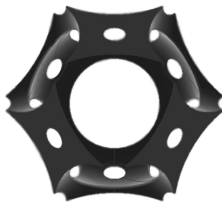
Objectives of micro-macro approaches :

- to link the microstructure to macroscopic parameters,
- to understand the material response according to various physics,
- to communicate between the different actors of the manufacturing.

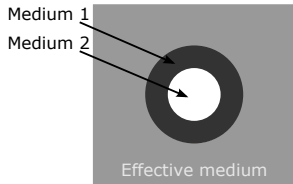
Approaches



a) Real microstructure

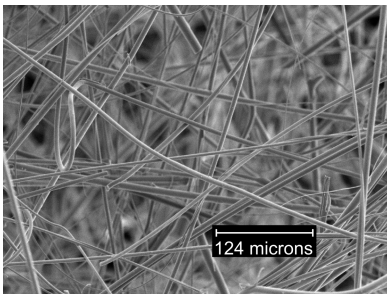
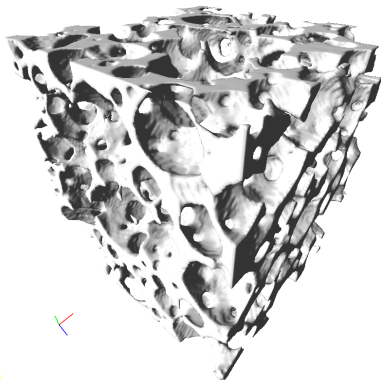


b) Idealized cell



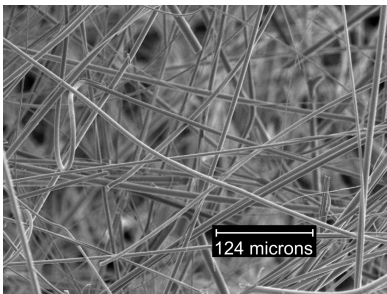
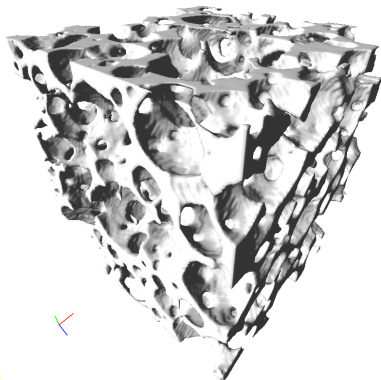
c) Self consistent model

Real microstructure



- + Real microstructure
- Expensive in memory and time
- Not adapted for optimisation

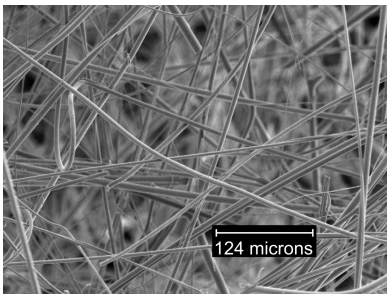
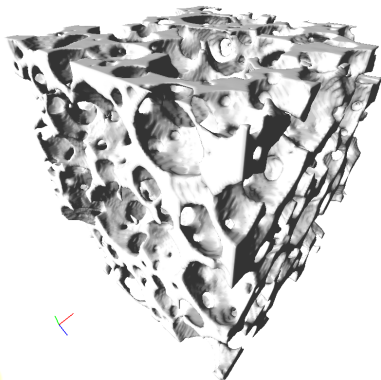
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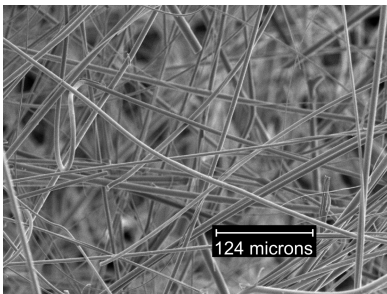
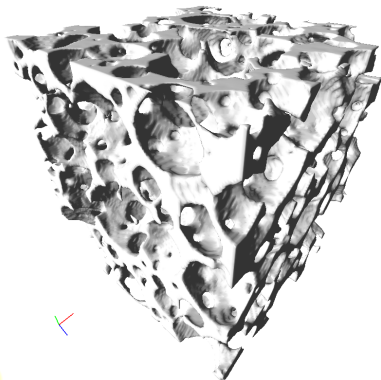
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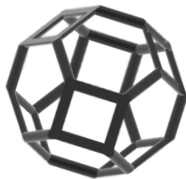
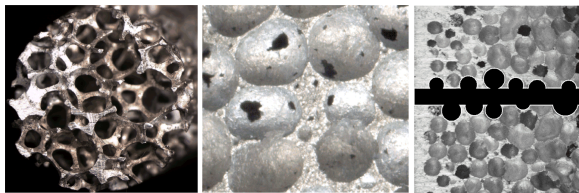
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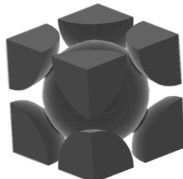


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Idealized cell



(a)



(b)

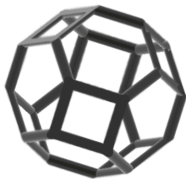
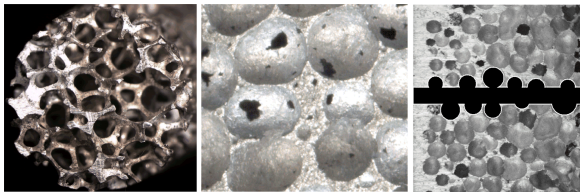


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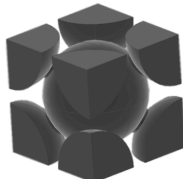
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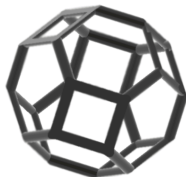
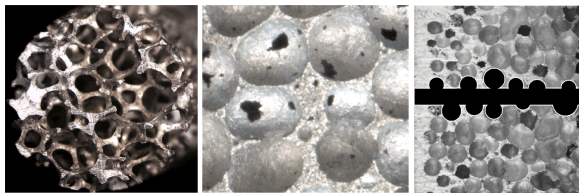


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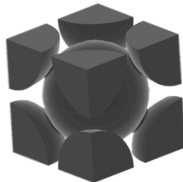
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Direct method: microscopic scale

- **Resolution** of two equations at the **microscopic scale** in dynamic regime
- Viscous effects (Navier-Stokes equation):

$$j\omega\rho_0\vec{u} = -\vec{\nabla}\tilde{p} + \eta\nabla^2\vec{u}$$

- Thermal effects (Heat conduction equation):

$$j\omega\rho_0c_p\tilde{\tau} = j\omega\tilde{p} + \kappa\nabla^2\tilde{\tau}$$

\vec{u} : velocity field

\tilde{p} : pressure field

$\tilde{\tau}$: temperature field

$\rho_0, c_p, \eta, \kappa$: properties of the fluid (air)

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- **Field integration**

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$$\tilde{\rho}_{eq}(\omega) = \frac{\eta}{j\omega \tilde{k}(\omega)}$$

- Thermal effects:

$$\phi \langle \tilde{\tau} \rangle = j\omega \frac{\tilde{k}'}{\kappa} \langle \tilde{p} \rangle$$

$$\tilde{\beta}(\omega) = \gamma - (\gamma - 1) \frac{j\omega \tilde{k}'(\omega)}{\nu' \phi} \quad \text{and} \quad \tilde{K}_{eq}(\omega) = \frac{\gamma P_0}{\phi \tilde{\beta}(\omega)}$$

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Hybrid method

The **hybrid method** contains two stages :

- Computation of the macroscopic parameters ($\phi, \sigma, \Lambda, \Lambda', \alpha_\infty, \dots$) from asymptotic behaviour (LF & HF) of the visco-inertial and thermal effects.
- Computation of the **dynamic mass density** $\tilde{\rho}(\omega)$ and of the **dynamic bulk modulus** $\tilde{K}(\omega)$ from semi-phenomenological models (JCA, JCAL, JCAPL, ...).

The two **main advantages** of this method are :

- the prediction on broad frequency range from only **three computations**,
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Hybrid method : summary

- **Geometric mesh :**

- porosity ϕ
- thermal characteristic length Λ'

- **LF viscous computation (Stokes):**

- static air flow resistivity σ
- static viscous tortuosity α_0 (Optional)

- **HF inertial computation (Perfect fluid)(\approx electrical conduction):**

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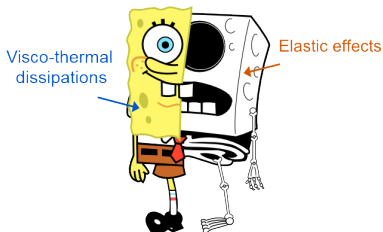
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Adding elastic effects



- Using Biot's theory, it is possible to couple elastic effects to any visco-thermal model

F.-X. Bécot, L. Jaouen, "Alternative Biot's form for porous media" (2013)

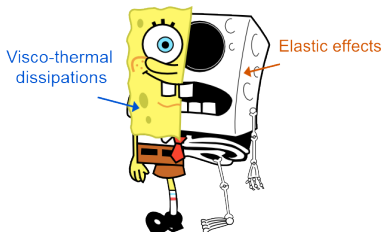
- Mixing laws: $E = E_1(1 - \phi_{meso}) + E_2\phi_{meso}$

J. C. Halpin , J. L. Kardos , "The Halpin-Tsai equations: A review" (1976)

- Micro-macro approaches for elastic properties

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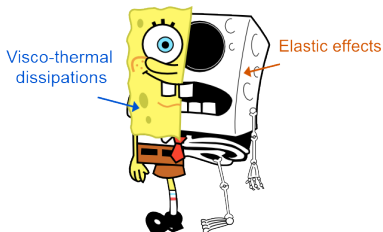
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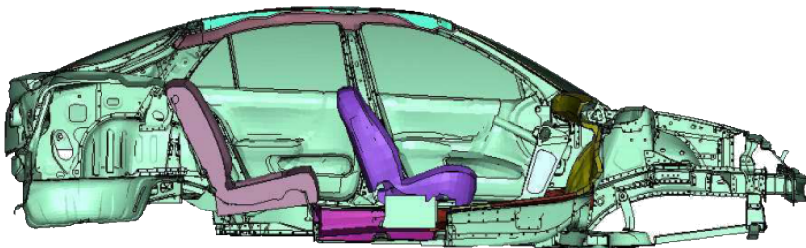
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Condensed models

- Example of a car with acoustical sound package

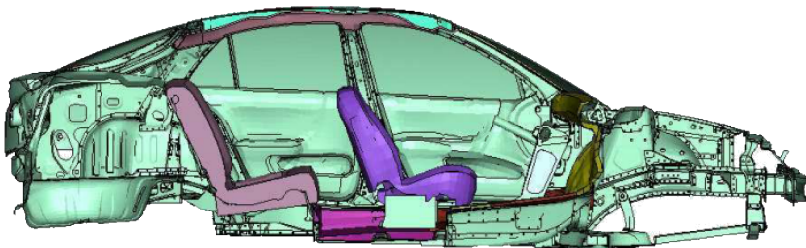


Source: A. Duval - Sapem 2005

- How to add physics to the models to reduce their size?

Condensed models

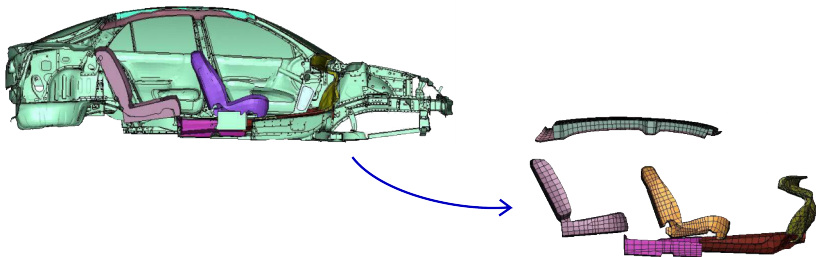
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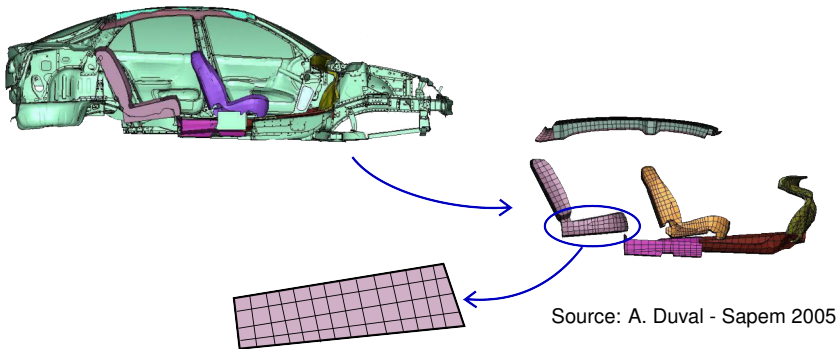
- How to add physics to the models to reduce their size?

- Focus on the porous parts



Source: A. Duval - Sapem 2005

- Focus on one porous part



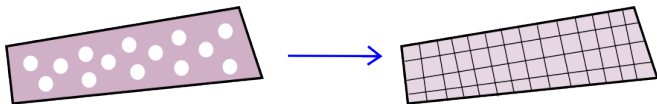
- Is it possible to model rigid or elastic meta-materials as a condensed model?

Non-conventional phenomena involved in meta-materials

- **Multi-scale material with diffusion processes**
 - Double porosity media (air inclusions)
 - Solid inclusions
 - Porous inclusions
 - Sorption
- **Acoustical resonators**
- **Soft membranes**
- **Periodicity**
- **Impedance matching**

Double porosity media

- Porous media with **air inclusions**



- Weight reduction
- Can take advantage from additional diffusion effects

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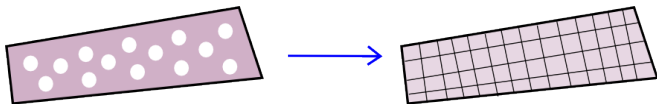
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Double porosity media

- Porous media with **air inclusions**

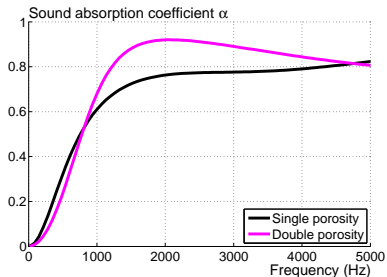


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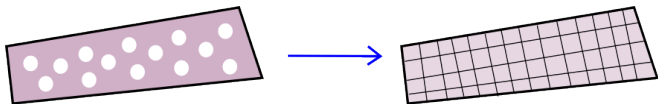
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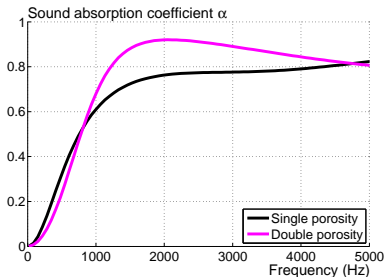


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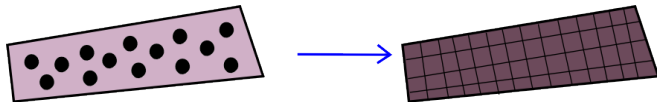


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Solid inclusions

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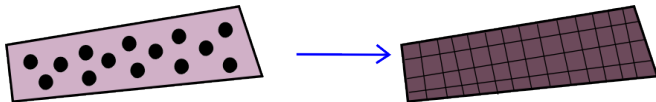
- Tortuous effect
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V. Tournat, V. Pagneux, D. Lafarge, L. Jaouen, "Multiple scattering of acoustic waves and porous absorbing media", *Phys. Rev. E* 70, 026609, (2004)

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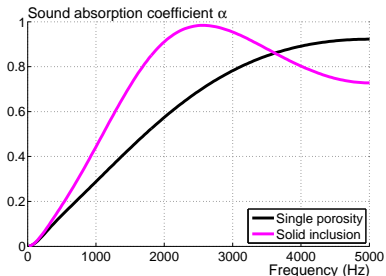
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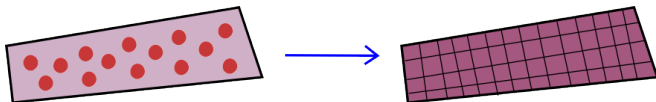
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Porous inclusions

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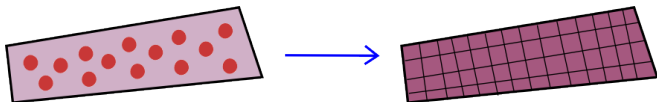


- Tortuous effect
- Permeability contrast
- Potential pressure diffusion effect

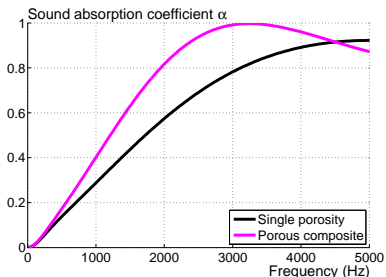
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Sorption

- Double porosity with nanoscale (third scale of porosity)
- The diffusion process can be enhanced by using sorption process (adsorption/desorption)
- Activated carbon for instance
- Reducing the bulk modulus (virtually increasing the volume of a cavity)

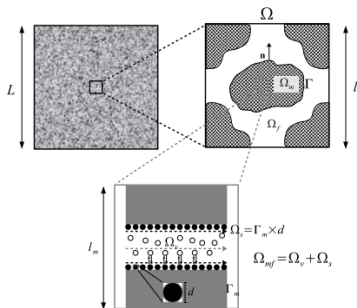


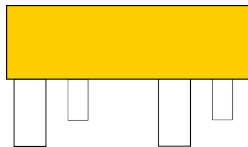
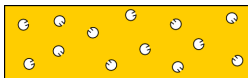
Fig. 1. Diagram of the scales of a hierarchical sorptive porous material.

R. Venegas, O. Umnova, "Influence of sorption on sound propagation in granular activated carbon", *J. Acoust. Soc. Am.* **68**, 162-181 (2017)

R. Venegas, C. Boutin, "Acoustics of sorptive porous materials.", *Wavemotion* **828**, 135-174 (2017)

Acoustical resonators

- Porous media with **(inner) resonators**



- Tortuous effect
- Helmholtz or quarter-wave resonator

$$\tilde{K}_{eq} = \frac{1}{\frac{(1 - \phi_{meso})}{\tilde{K}_{eq_micro}} + \frac{\phi_{meso}}{\tilde{K}_{eq_res}}}$$

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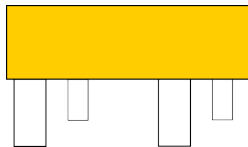
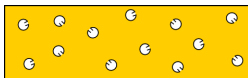
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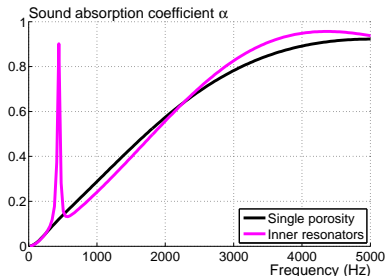
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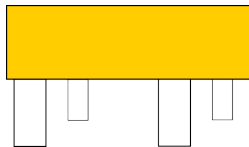
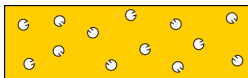
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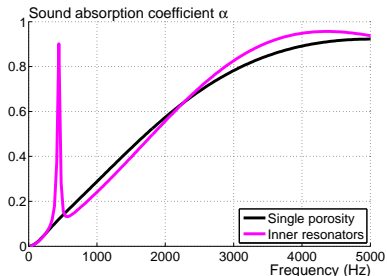
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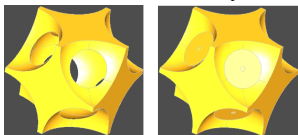
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Soft membranes

- Rigid membranes are known to increase the airflow resistivity as well as the tortuosity



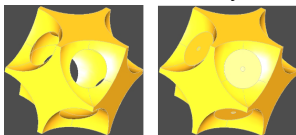
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C. Gaulon, J. Pierre, C. Derec, L. Jaouen, F.-X. Bécot, F. Chevillotte, F. Elias, W. Drenckhan, and V. Leroy, "Acoustic absorption of solid foams with thin membranes.", *Appl. Phys. Lett.* 112(261904), (2018)

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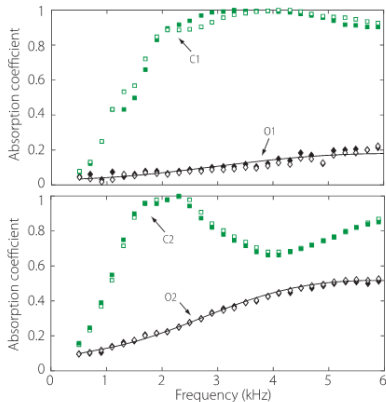


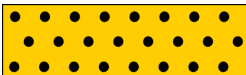
FIG. 3. Measured absorption versus frequency for open cell foams (O1 and O2, black symbols) and closed-cell foams (C1 and C2, green symbols). Samples were 2 cm thick. Reproducibility was tested by measuring 2 samples of each type (solid and open symbols). Solid black lines show JCAL model for the open-cell foams.

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Periodicity

- Porous media with **periodic** solid inclusions

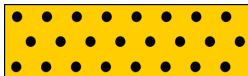


- Tortuous effect
- Bragg interferences

- The waves are reflected, not dissipated!

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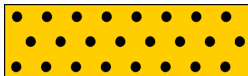


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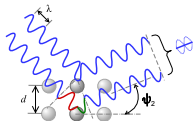
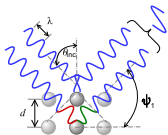
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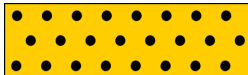
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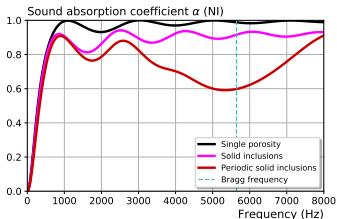
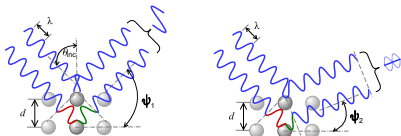
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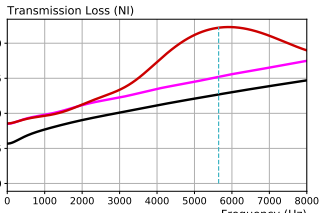
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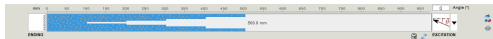


Impedance matching

- The **impedance matching** principle consists in **gradually modify the impedance** to tend towards **zero reflection**
- Generally achieved by **modifying the topology**
- Functionnally **graded materials** are also good candidates

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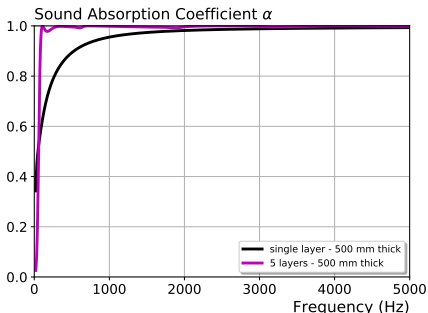
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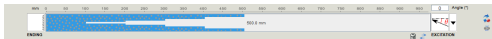


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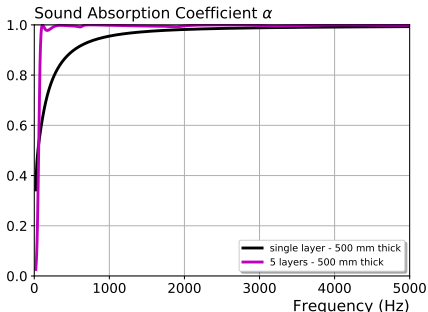
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Non-conventional phenomena involved in meta-materials

Summary of analytical condensed models:

- **Multi-scale material with diffusion processes**

- Double porosity media (air inclusions) 
- Solid inclusions 
- Porous inclusions 
- Sorption 

- **Acoustical resonators** 

- **Soft membranes** 

- **Periodicity**





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How to condense any non-conventional effect?

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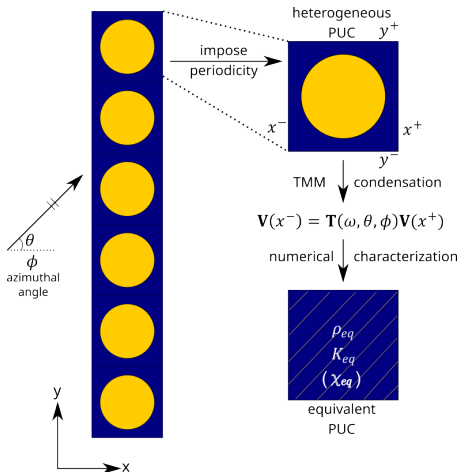
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How to condense any non-conventional effect?

Numerical condensation



- A. Parrinello, A., G. Ghiringhelli "Transfer matrix representation for periodic planar media," Journal of Sound and Vibration 371, 196–20, 2016
- A. Parrinello, A., G. Ghiringhelli, N. Atalla "Generalized transfer matrix method for periodic planar media," Journal of Sound and Vibration 464, 11499, 2020
- A. Sreekumar, F. Chevillotte, E. Gourdon "Numerical characterization of heterogeneous meta-materials." Forum Acusticum 2023, Turin, Italy. Submitted to J. Acous. Soc. Am.

Numerical condensation

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$$\mathbf{T} = \begin{bmatrix} \cos(k_x l) & j \frac{\rho_{\text{eq}} \omega}{k_x} \sin(k_x l) \\ j \frac{k_x}{\rho_{\text{eq}} \omega} \sin(k_x l) & \cos(k_x l) \end{bmatrix} \quad \text{with } k_x = \sqrt{k_{\text{eq}}^2 - k_t^2}$$

and $k_t = k_0 \sin(\theta)$

- Using the Stroh formalism:

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χ enables to account for asymmetry

J.-P. Groby, M. Malléjac, A. Merkel, V. Romero-Garcia, V. Tournat, D. Torrent, and J. Li, "Analytical modeling of one-dimensional resonant asymmetric and reciprocal acoustic structures as Willis materials," *New Journal of Physics* 23(5), 466 053020, 2021

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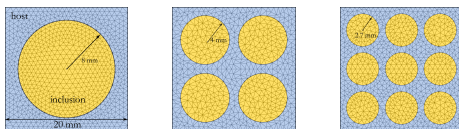
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Numerical condensation

- Configurations and material parameters



	Host	Inclusion				
		Air	Porous I	Porous II	Porous III	Rigid
σ [N · m · s ⁻⁴]	8900	-	2×10^4	2×10^5	2×10^7	-
ϕ	0.95	-	0.95	0.95	0.95	-
α_∞	1.42	-	1	1	1	-
Λ [m]	100×10^{-6}	-	8.802×10^{-5}	2.783×10^{-5}	2.783×10^{-6}	
Λ' [m]	360×10^{-6}	-	8.802×10^{-5}	2.783×10^{-5}	2.783×10^{-6}	

Numerical condensation

- Sound absorption coefficients under DF (1/3)



A

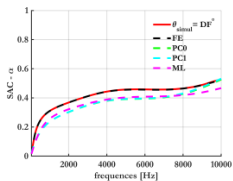
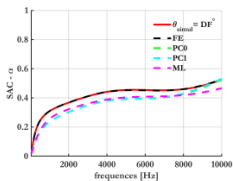
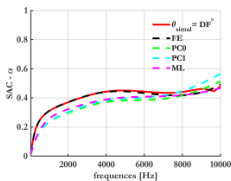


B

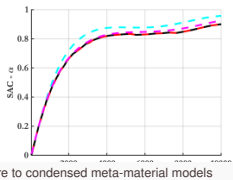
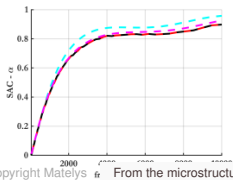
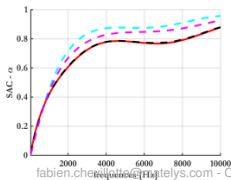


C

Air

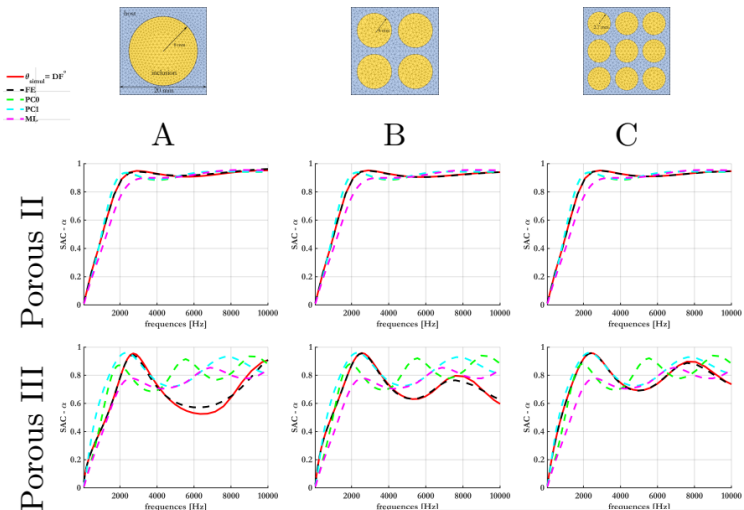


Porous I



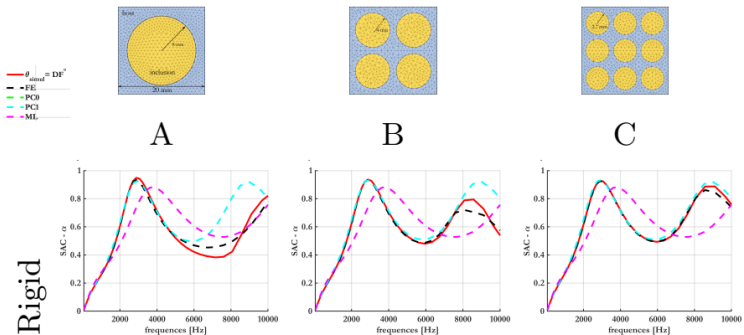
Numerical condensation

- Sound absorption coefficients under DF (2/3)



Numerical condensation

- Sound absorption coefficients under DF (3/3)



Numerical condensation

- Pressure fields @ 3064 Hz (1/3)



A

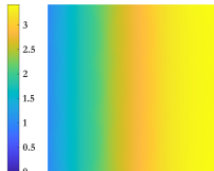
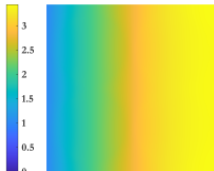


B

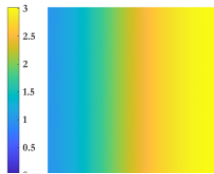
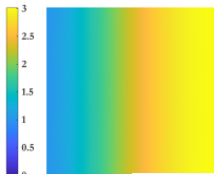


C

Air



Porous I



Numerical condensation

- Pressure fields @ 3064 Hz (2/3)

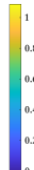
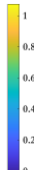
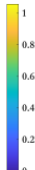
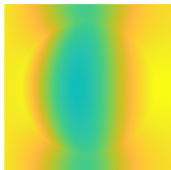


A

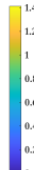
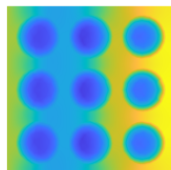
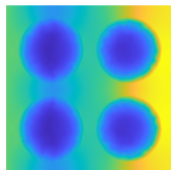
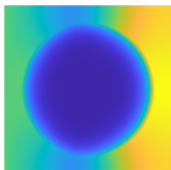
B

C

Porous II



Porous III



Numerical condensation

- Pressure fields @ 3064 Hz (3/3)



A

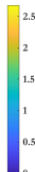
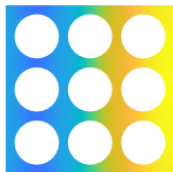
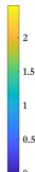
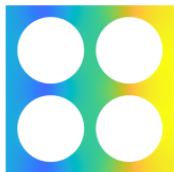
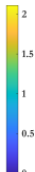


B



C

Rigid



Conclusion

- The equivalent fluid formalism enables to embed complex phenomena into condensed models (flow distortion, pressure diffusion, inner resonance, ...).
- Simple micro- or meso-structures can be described from analytical formulas
- Micro-macro methods can be used when dealing with more complex shapes
- Numerous homogenisation models are available for non-conventional phenomena involved in meta-materials.
- Otherwise, numerical characterization procedure can be employed.

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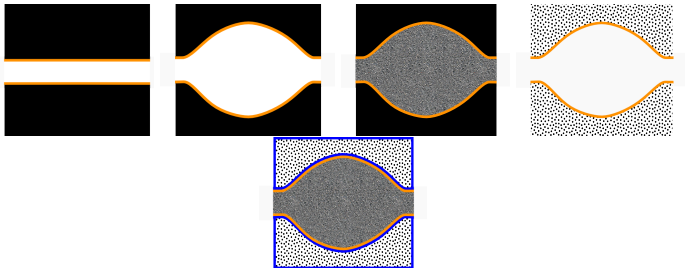
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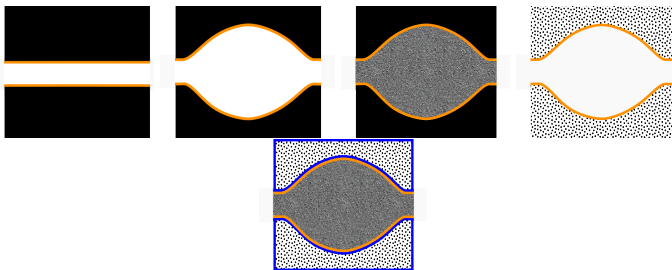
Composite model

- A composite model has recently been presented.
F. Chevillotte, L. Jauouen, F.-X. Bécot, "On the modeling of visco-thermal dissipations in heterogeneous porous media", *J. Acoust. Soc. Am.* **138**(6), 3922-3929 (2015)
- Consideration of the **shape** of the mesoscopic parts.
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Perforated plate

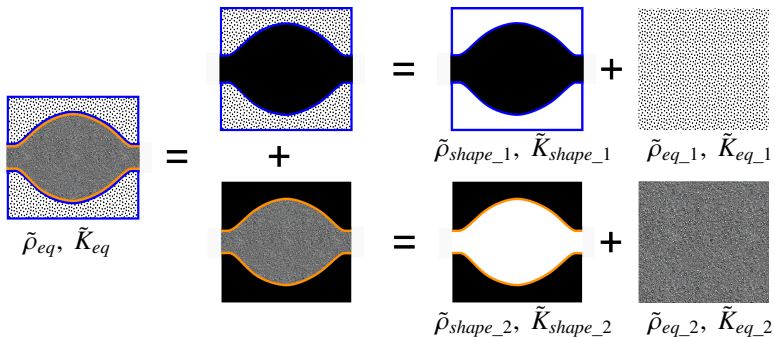
Usual model

Porous inclusion

Double porosity

Composite

Composite porous media



Shape consideration

Visco-inertial effects

$$\tilde{\rho}_{shape}(\omega) = \tilde{\rho}_{eq}(\omega)\alpha_{\infty} \left[1 - j\frac{\omega_v}{\omega} \tilde{G}(\omega) \right]$$

$$\tilde{G}(\omega) = \sqrt{1 + \frac{1}{2}jM\frac{\omega}{\omega_v}}$$

$$M = \frac{8k_0\alpha_{\infty}}{\phi\Lambda^2}; \quad \omega_v = \frac{\nu\phi}{k_0\alpha_{\infty}}$$

$$\nu = \frac{\eta}{\tilde{\rho}_{eq}(\omega)}$$

Thermal effects

$$\tilde{K}_{shape}(\omega) = \frac{\tilde{K}_{eq}(\omega)}{\gamma - (\gamma - 1) \left[1 - j\frac{\omega_t}{\omega} \tilde{G}'(\omega) \right]^{-1}}$$

$$\tilde{G}'(\omega) = \sqrt{1 + \frac{1}{2}jM'\frac{\omega}{\omega_t}}$$

$$M' = \frac{8k'_0}{\phi\Lambda'^2}; \quad \omega_t = \frac{\nu'\phi}{k'_0}$$

$$\nu' = \frac{\kappa}{\tilde{\rho}_{eq}(\omega)C_p}$$

Porous composite

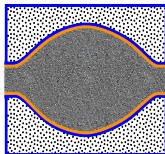
$$\tilde{\rho}_{eq} = \frac{1}{\frac{1}{\tilde{\rho}_{eq_shape_1}} + \frac{1}{\tilde{\rho}_{eq_shape_2}}}$$

$$\tilde{K}_{eq} = \frac{1}{\frac{\tilde{F}_{dw}^1}{\tilde{K}_{eq_shape_1}} + \frac{\tilde{F}_{dw}^2}{\tilde{K}_{eq_shape_2}}}$$

$$\phi_{meso} = \phi_{shape_1}$$

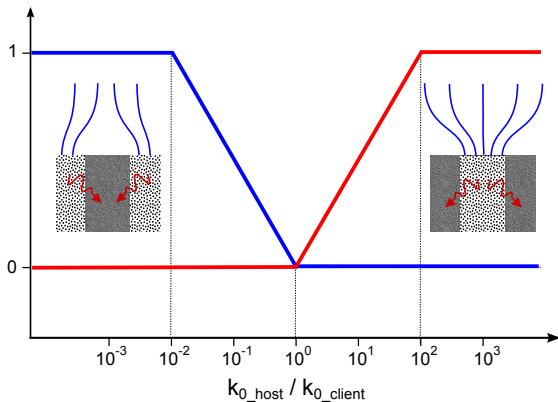
$$1 - \phi_{meso} = \phi_{shape_2}$$

\tilde{F}_{dw}^1 et \tilde{F}_{dw}^2 : Diffusion functions 2 → 1 et 1 → 2.

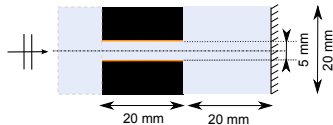
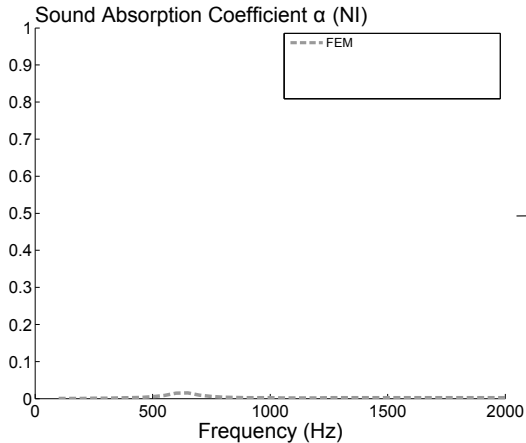


Consideration of diffusion and flow distortion

Weighting functions

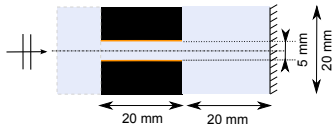
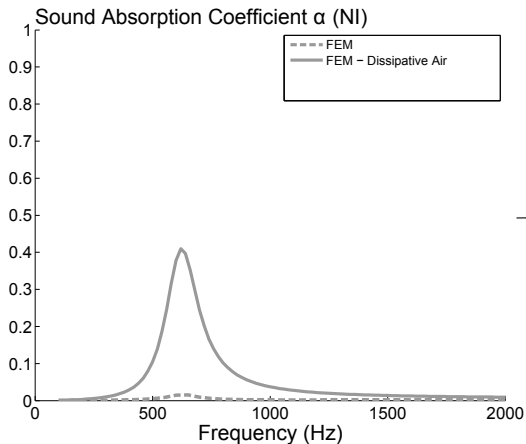


Perforated plate



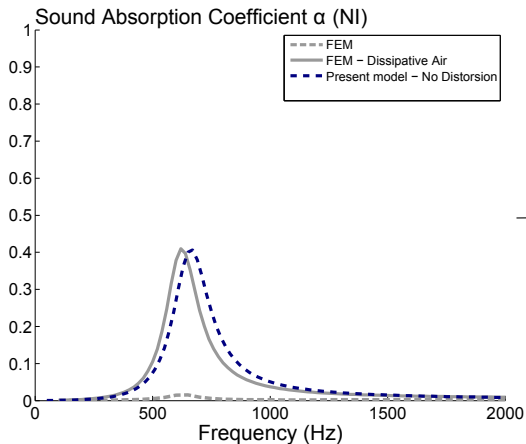
- Validation with flow distortion.

Perforated plate

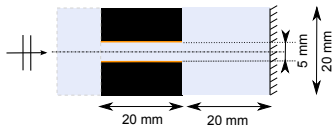


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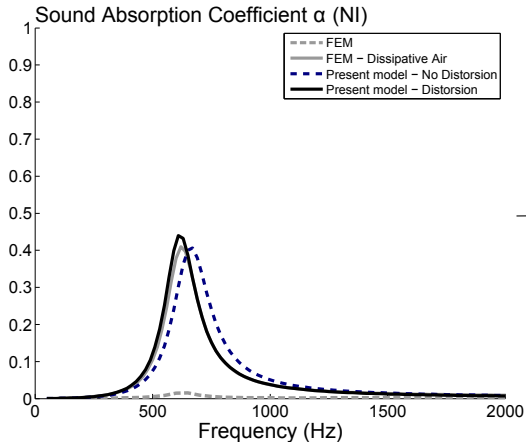
Perforated plate



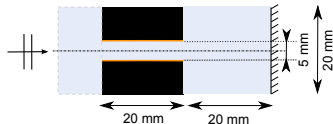
- Validation with flow distorsion.



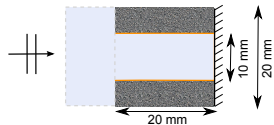
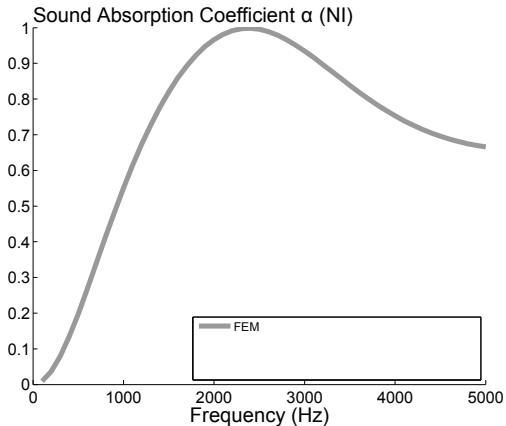
Perforated plate



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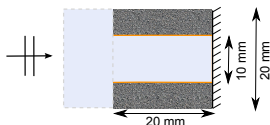
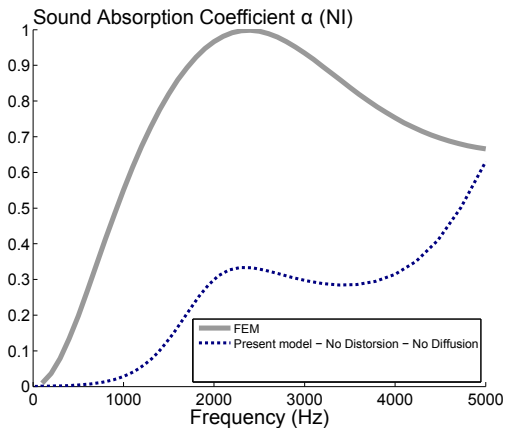
Double porosity media



Host: $\sigma = 1\,000\,000 \text{ N.s.m}^{-4}$
Client: Air

- Validation with diffusion.

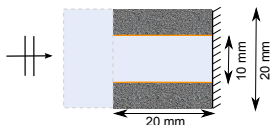
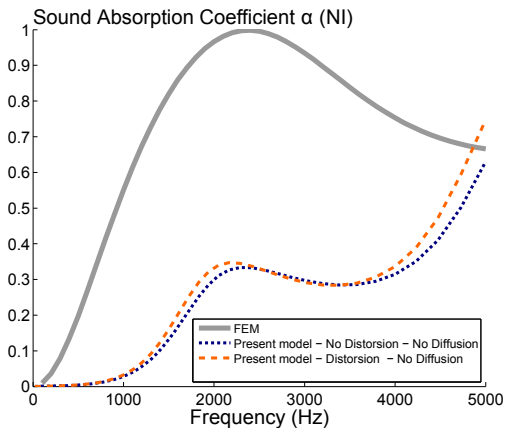
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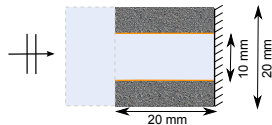
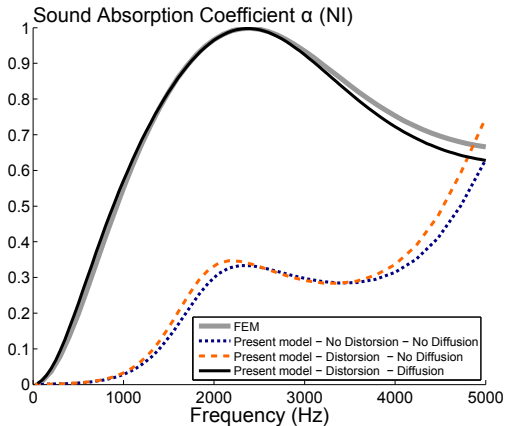
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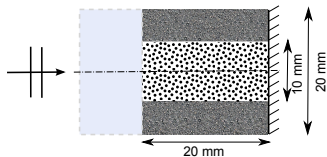
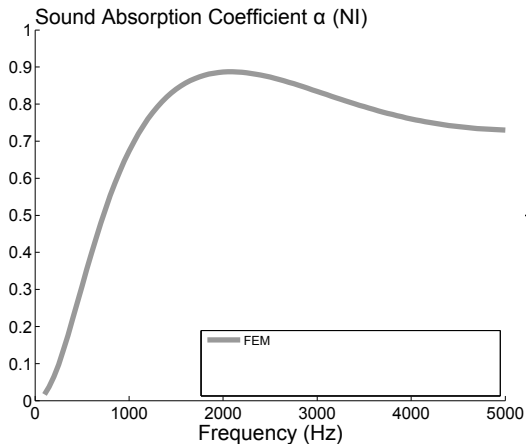
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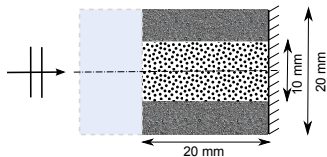
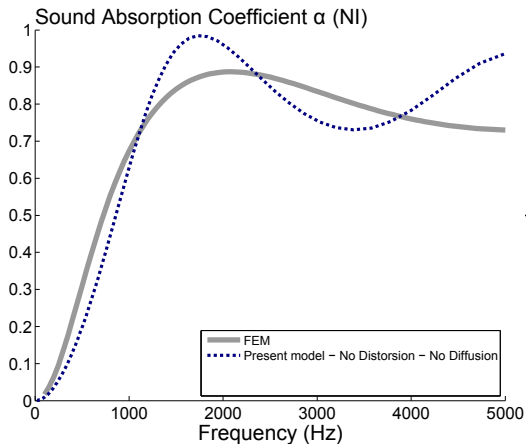
Porous composite - simple shape



Host: $\sigma = 1\,000\,000 \text{ N.s.m}^{-4}$
Client: $\sigma = 11\,500 \text{ N.s.m}^{-4}$

- Validation with diffusion between porous media.

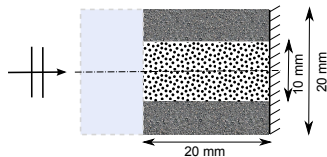
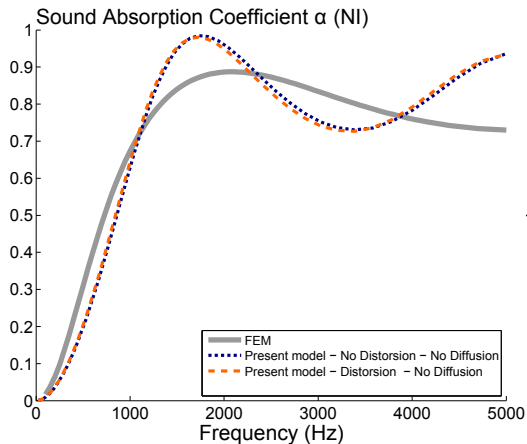
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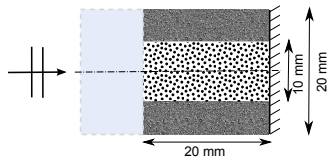
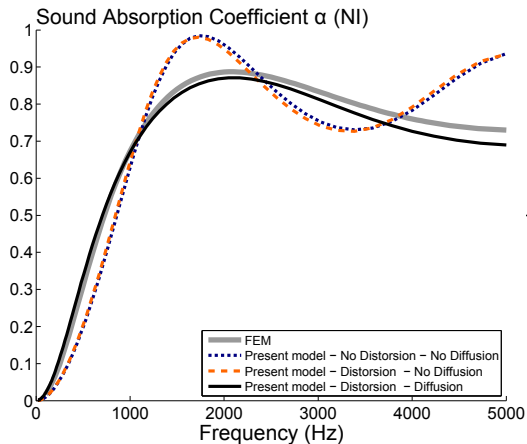
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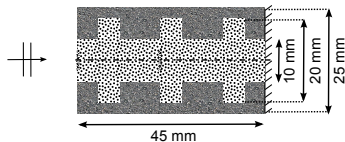
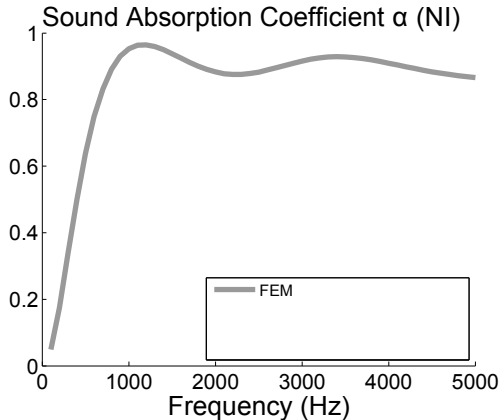
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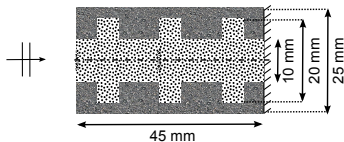
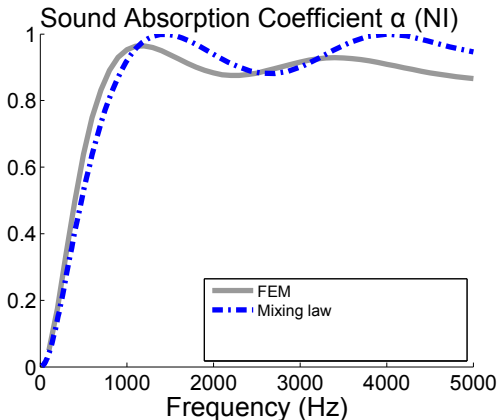
Porous composite - complex shape



Host: $\sigma = 1\,000\,000 \text{ N.s.m}^{-4}$
Client: $\sigma = 11\,500 \text{ N.s.m}^{-4}$

- Validation with shape and diffusion.

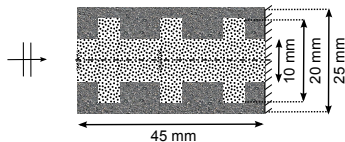
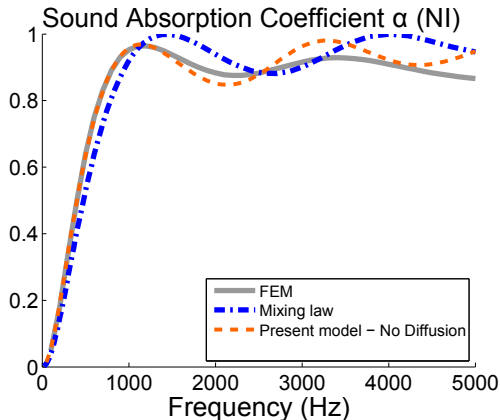
Porous composite - complex shape



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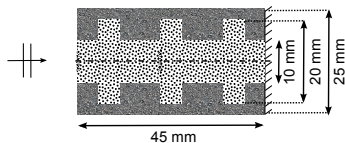
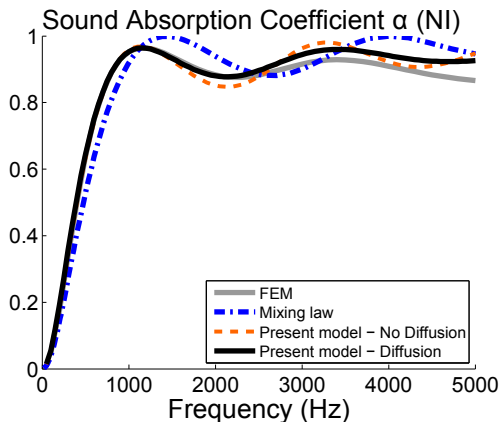
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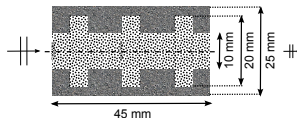
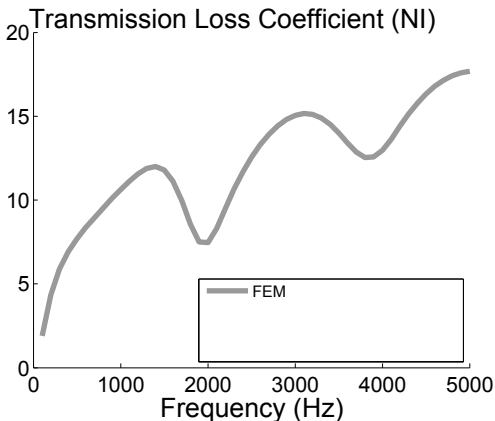
Porous composite - complex shape



Host: $\sigma = 1\,000\,000 \text{ N.s.m}^{-4}$
Client: $\sigma = 11\,500 \text{ N.s.m}^{-4}$

- Validation with shape and diffusion.

Porous composite - Transmission



Host: $\sigma = 150\,000 \text{ N.s.m}^{-4}$

$E = 500\,000 \text{ Pa}$

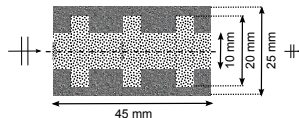
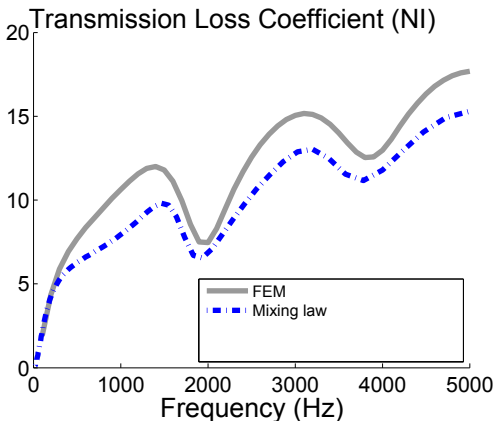
Client: $\sigma = 10\,000 \text{ N.s.m}^{-4}$

$E = 200\,000 \text{ Pa}$

- Consideration for elastic effects:

J. C. Halpin , J. L. Kardos , "The Halpin-Tsai equations: A review", Polymer Engineering and Science, 16 (5), 344 - 352 1976).

Porous composite - Transmission



Host: $\sigma = 150\,000 \text{ N.s.m}^{-4}$

$E = 500\,000 \text{ Pa}$

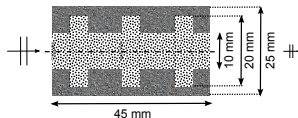
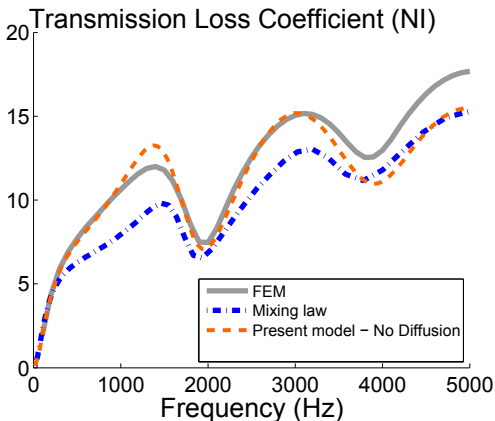
Client: $\sigma = 10\,000 \text{ N.s.m}^{-4}$

$E = 200\,000 \text{ Pa}$

- Consideration for elastic effects:

J. C. Halpin , J. L. Kardos , "The Halpin-Tsai equations: A review", Polymer Engineering and Science, 16 (5), 344 - 352 1976).

Porous composite - Transmission



Host: $\sigma = 150\,000 \text{ N.s.m}^{-4}$

$E = 500\,000 \text{ Pa}$

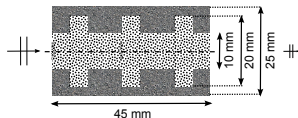
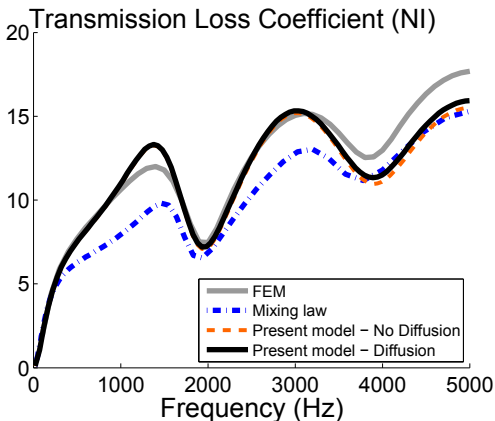
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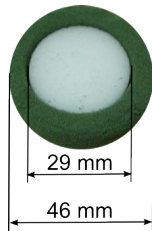
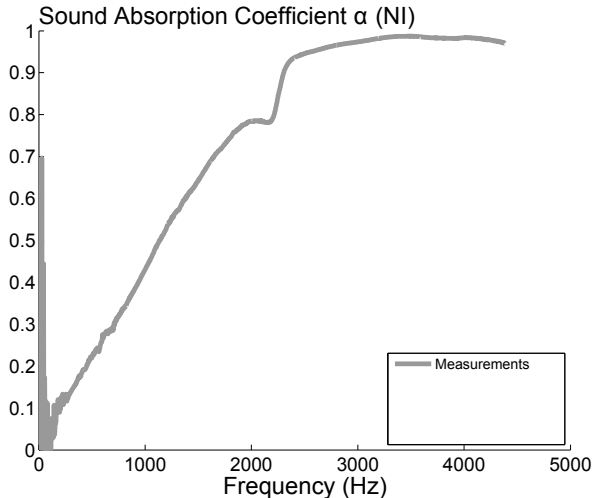
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Experimental validation



Normal incidence

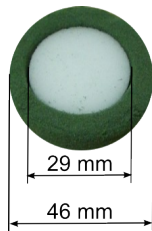
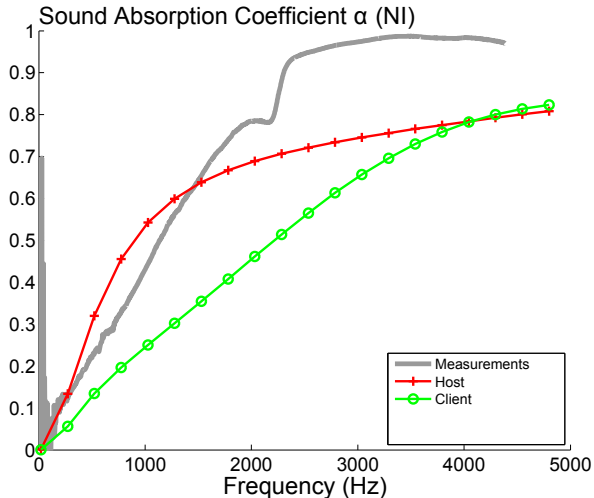
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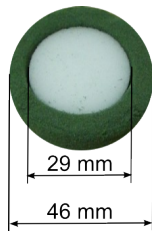
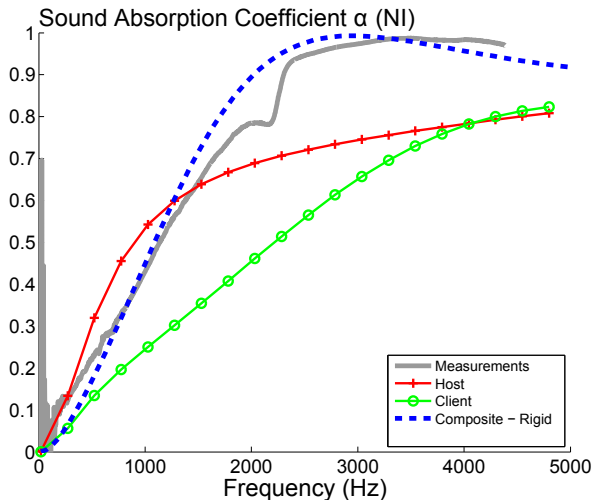
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Normal incidence

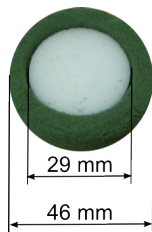
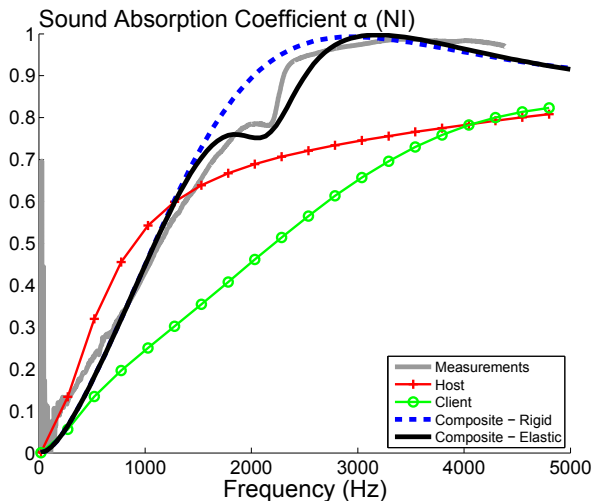
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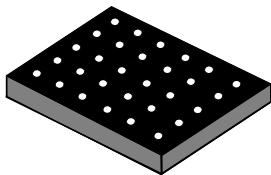
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Client: $\sigma = 10\,000 \text{ N.s.m}^{-4}$

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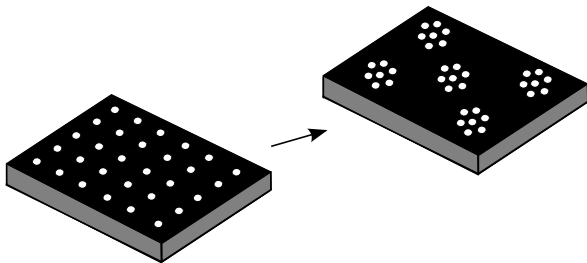
Examples with shape consideration

- Modified perforated plates.



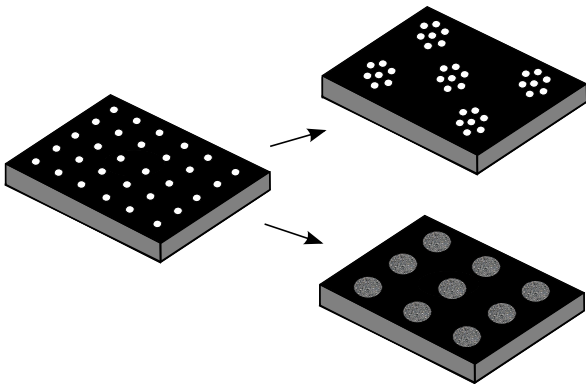
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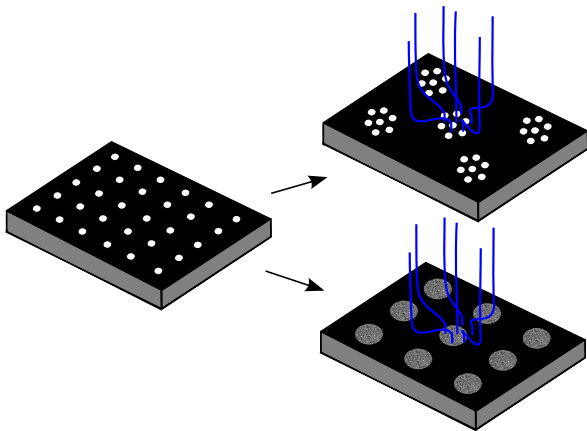
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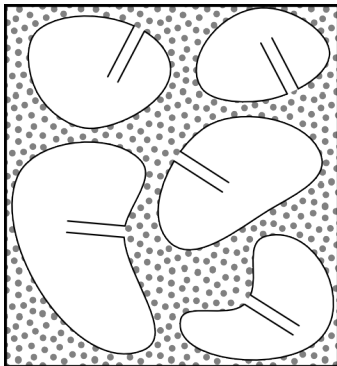
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Examples with shape consideration

- Inner resonators.



C. Boutin, "Acoustics of rigid porous media with inner resonators", J. Acoust. Soc. Am. 134 , 4717 (2013).

Alternative concepts

- **Mechanical resonators** can also be added to the structure
- **Bending mode resonators**

- **Spring-mass resonators**

F. Tateo, J. Michiels, I. Lopez Arteaga, and H. Nijmeijer. "Resonant lattices for low-frequency vibro-acoustic control." In Novem, Dubrovnik, Croatia, 2015.

C. Claeys, E. Deckers, B. Pluyfymers, and W. Desmet. "A lightweight vibro-acoustic metamaterial demonstrator: numerical and experimental investigation." In Novem, Dubrovnik, Croatia, 2015.

- **Non-linear mechanical resonators**

F. Vakakis and O. Gendelman. "Energy pumping in nonlinear mechanical oscillators: Part II -resonance capture." ASME. J. Appl. Mech., 68(1): 42-48, 2000.

E. Gourdon, N. A. Alexander, C. A. Taylor, C. H. Lamarque, and S. Pernot. "Nonlinear energy pumping under transient forcing with strongly nonlinear coupling: Theoretical and experimental results." J. Sound Vib., 300(3-5): 522-551, 2007.

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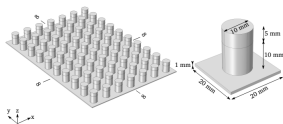


Figure 2. Considered elastic waveguide: 2D periodic lattice (left); unit cell (right).

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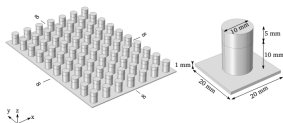


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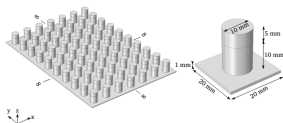


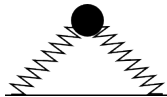
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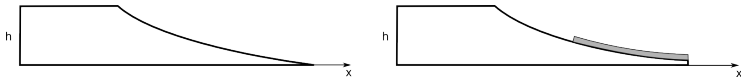


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- **Impedance matching concept applied to structural vibrations (« Acoustic Black Hole »)**



C. L. Pekeris. "Theory of Propagation of Sound in a Half-Space of Variable Sound Velocity under Conditions of Formation of a Shadow Zone." J. Acoust. Soc. Am., 18 :295–315, 1946.

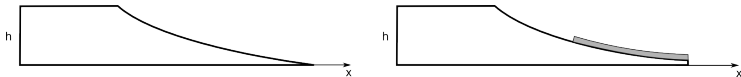
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- **Meta-material for underwater acoustics are mainly based on solid visco-elastic matrix with soft inclusions)**

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