

# Strain wave propagation through felt-like material

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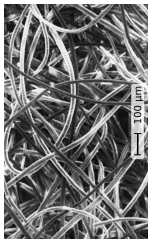
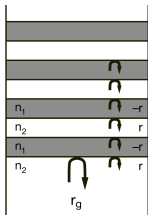


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# Motivation



- Metamaterials are engineered materials<sup>1</sup> designed to have properties that can not be found in natural materials
- They are created by arranging or structuring components at a smaller scale than the wavelength of the phenomena they interact with
- Metamaterials can be designed such that they exhibit band gaps (BG) and properties like negative group velocity (NGV)<sup>2,3</sup>
- These materials can be used to manipulate acoustic waves and, therefore beneficial for noise and vibration reduction and control

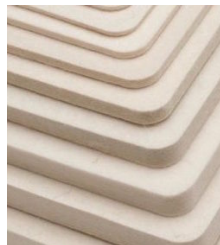
<sup>1</sup>S. Guenneau and R. V. Craster, *Fundamentals of Acoustic Metamaterials* Dordrecht: Springer Netherlands, 2013, pp. 1–42.

<sup>2</sup>N. Fang, D. Xi, J. Xu, M. Ambati, W. Srituravanich, C. Sun, X. Zhang, “Ultrasonic metamaterials with negative modulus,” *Nat. Mater.*, **5**(6), pp. 452–456, 2006.

<sup>3</sup>S. Guenneau, A. Movchan, G. Petursson, S. A. Ramakrishna, “Acoustic metamaterials for sound focusing and confinement,” *New J. Phys.*, **9**(11), p. 399, 2007.

# Main aim

- Numerically analyse the acoustic wave propagation through a felt continuum using a previously obtained visco-elastic model equation is performed
- Dispersion and dissipation analysis is presented
- Studying numerically the possible effect of NGV and BG on the wave propagation



# Model equation

A 1D equation for modelling wave propagation through a felt-like medium was derived in <sup>4</sup>. Presented here in dimensionless variables in the following form:

$$\frac{\partial^2 (\varepsilon^p)}{\partial x^2} - \frac{\partial^2 \varepsilon}{\partial t^2} + \frac{\partial^3 (\varepsilon^p)}{\partial x^2 \partial t} - \delta \frac{\partial^3 \varepsilon}{\partial t^3} = 0, \quad (1)$$

where  $\varepsilon(x, t)$  is the strain,  $p \geq 1$  is the material parameter and  $0 < \delta \leq 1$  is the hereditary parameter.

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<sup>4</sup>D. Kartofelev, A. Stulov, "Wave propagation and dispersion in microstructured wool felt," *Wave Motion*, **57**, pp. 23–33, 2015.

# Dispersion and dissipation analysis

Linearised form of Eq. (1) is obtained by taking  $p = 1$

$$\frac{\partial^2 \varepsilon}{\partial x^2} - \frac{\partial^2 \varepsilon}{\partial t^2} + \frac{\partial^3 \varepsilon}{\partial x^2 \partial t} - \delta \frac{\partial^3 \varepsilon}{\partial t^3} = 0. \quad (2)$$

The solution to (2) can be assumed in the form

$$\varepsilon(x, t) \propto \exp(\sigma t + ikx), \quad (3)$$

where  $\sigma$  is the attenuation,  $i$  is the imaginary unit and  $k$  is the wavenumber. Substituting (3) into (2) produces the following characteristic equation:

$$\delta \sigma^3 + \sigma^2 + k^2 \sigma + k^2 = 0. \quad (4)$$

Since  $k$  is considered to be real the  $\sigma$  can be written in terms of

$$\sigma(k) = \varsigma(k) + i\mu(k). \quad (5)$$

# Dispersion and dissipation analysis

Using (5) the assumption (3) takes the following form

$$\varepsilon(x, t) \propto \exp((\varsigma + i\mu)t + ikx) = \exp(\varsigma t) \cdot \exp(i(\mu t + kx)). \quad (6)$$

Where  $\varsigma$  acts in time as an attenuation coefficient and  $\mu$  is characterising the dispersion. The equation (4) can be solved using (5) as a system and written as a function of wavenumber  $k$ :

$$\begin{cases} \varsigma = \frac{1}{12\delta S} \left[ \sqrt[3]{4S^2 - 4S + 2\sqrt[3]{2}(1 - 3k^2\delta)} \right], \\ \mu = \frac{\sqrt{6}}{12\delta S} \sqrt{\sqrt[3]{2}S^4 - 4S^2(1 - 3k^2\delta) + 2\sqrt[3]{4}(1 - 3k^2\delta)^2}, \end{cases} \quad (7)$$

where

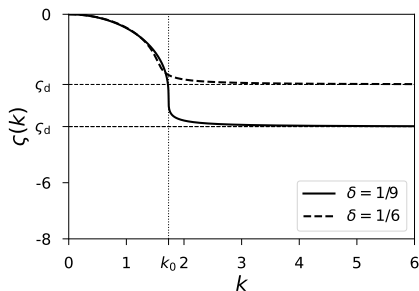
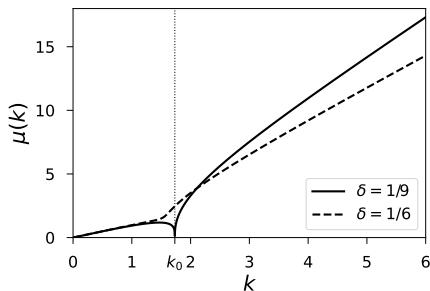
$$S = \sqrt[3]{2 - 9k^2\delta(1 - 3\delta) + 3k\delta\sqrt{3Q}}, \quad (8)$$

$$Q = 4k^4\delta - k^2(1 + 18\delta - 27\delta^2) + 4. \quad (9)$$

# Dispersion and dissipation analysis

Three different solution regimes are distinguishable depending on the value of  $\delta$ :

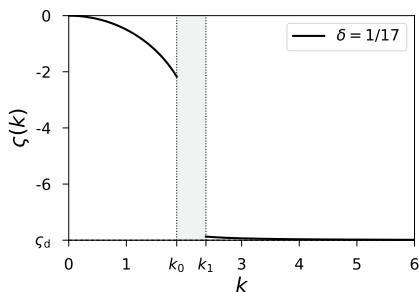
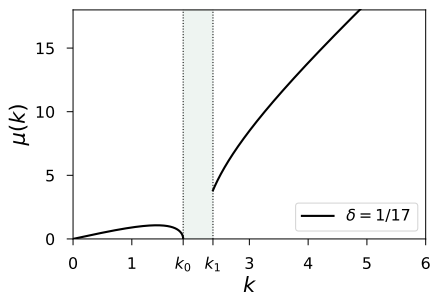
- 1 Non-dispersive and non-dissipative regime, for  $\delta = 1$
- 2 **Low hereditary** regime with a continuous dispersion and dissipation curves, for  $1/9 \leq \delta < 1$



# Dispersion and dissipation analysis

Three different solution regimes are distinguishable depending on the value of  $\delta$ :

- ③ **High hereditary** regime, when  $0 < \delta < 1/9$  dispersion and dissipation graphs have a BG region



At the limit  $\delta \rightarrow 0$  the BG widens since  $k_0 \rightarrow 2$  and  $k_1 \rightarrow \infty$ .

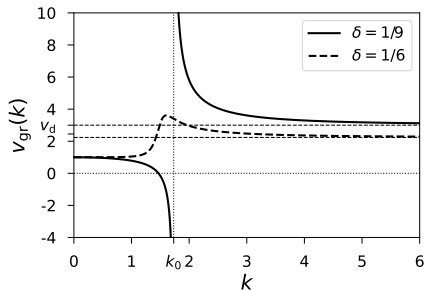


# Negative group velocity

The group velocity is given by

$$v_{\text{gr}}(k) = \frac{d\mu}{dk}. \quad (10)$$

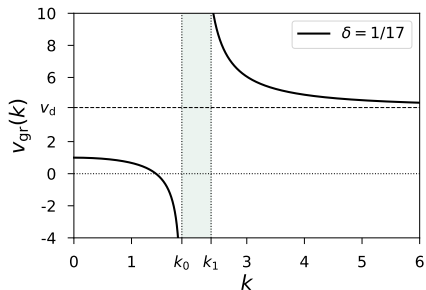
A region of NGV appears for  $0 < \delta \lesssim 0.134$ .



Here  $\delta = 1/9 \approx 0.111$ , and  
 $\delta = 1/6 \approx 0.167$

# Negative group velocity

For  $0 < \delta < 1/9$  the NGV region is neighboured by the BG.



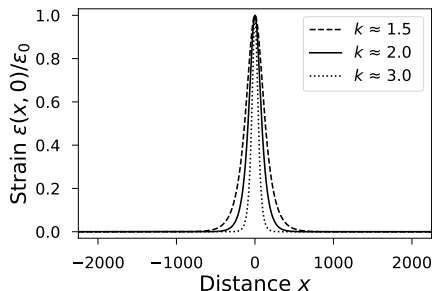
## The initial value problem (IVP), $\delta = 1/17$

$$\frac{\partial^2 \varepsilon}{\partial x^2} - \frac{\partial^2 \varepsilon}{\partial t^2} + \frac{\partial^3 \varepsilon}{\partial x^2 \partial t} - \delta \frac{\partial^3 \varepsilon}{\partial t^3} = 0$$

A bell-shaped initial condition (IC) is selected in the following form:

$$\varepsilon(x, 0) = \varepsilon_0 \operatorname{sech}(\beta x), \quad (11)$$

where the characteristic width and amplitude are  $\beta$  and  $\varepsilon_0$ , respectively.



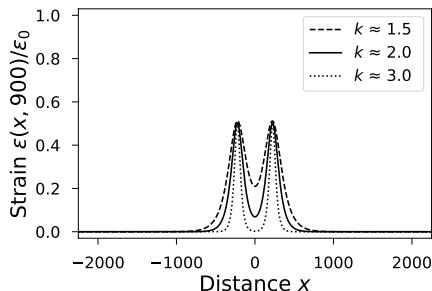
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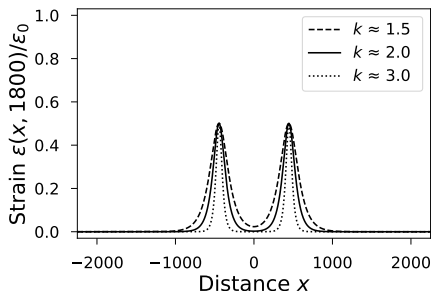
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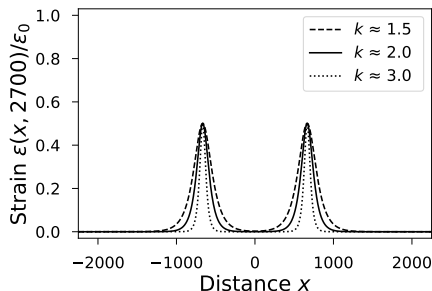
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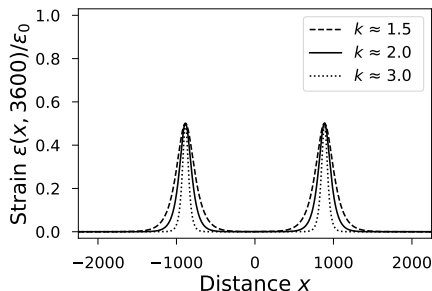
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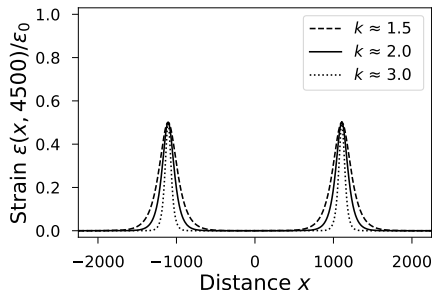
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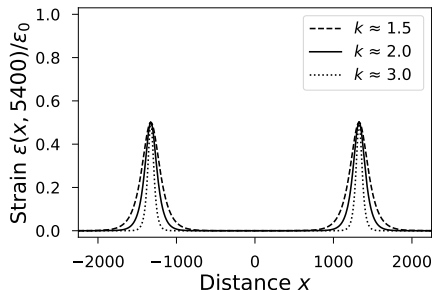
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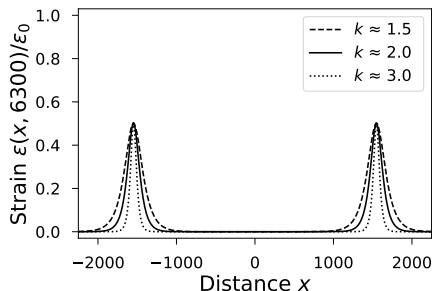
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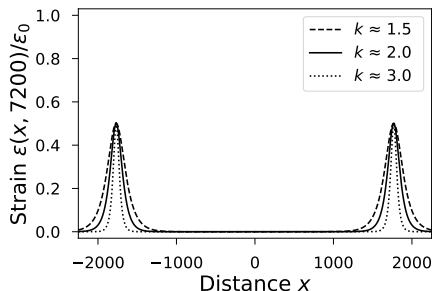
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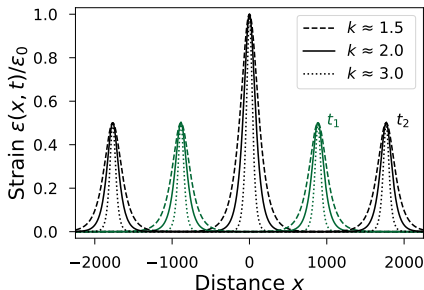
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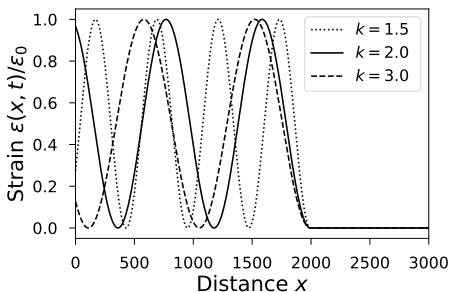
## Boundary value problem (BVP), $\delta = 1/17$

$$\frac{\partial^2 \varepsilon}{\partial x^2} - \frac{\partial^2 \varepsilon}{\partial t^2} + \frac{\partial^3 \varepsilon}{\partial x^2 \partial t} - \delta \frac{\partial^3 \varepsilon}{\partial t^3} = 0$$

A sinusoidal boundary condition (BC) in the form

$$\varepsilon(0, t) = \frac{\varepsilon_0}{2} [1 - \cos(2\pi\gamma t)], \quad (12)$$

where  $\varepsilon_0$  is the initial amplitude of the signal,  $\gamma$  is the signal frequency.



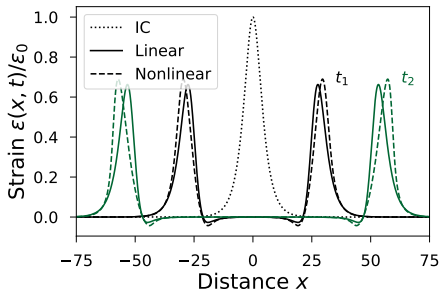
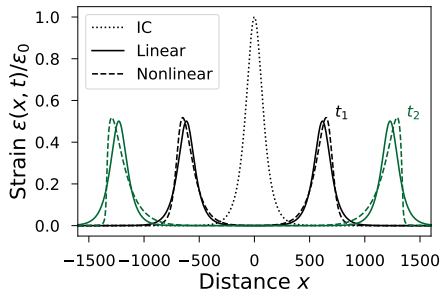
# Linear vs. weakly nonlinear, $\delta = 1/17$

$$\frac{\partial^2 (\varepsilon^p)}{\partial x^2} - \frac{\partial^2 \varepsilon}{\partial t^2} + \frac{\partial^3 (\varepsilon^p)}{\partial x^2 \partial t} - \delta \frac{\partial^3 \varepsilon}{\partial t^3} = 0$$

The IC is selected in the following form:

$$\varepsilon(x, 0) = \varepsilon_0 \operatorname{sech}(\beta x).$$

A linear case where  $p = 1$  and a weakly nonlinear case where  $p = 1.1$  for  $k \approx 2$  (left) and  $k \approx 50$  (right).



# Conclusions

- Model equation for modelling strain wave propagation through felt was analysed
- Three solution regimes were identified depending on the values of  $\delta$
- Dispersion and dissipation analysis of the model was presented
- The BG and NGV effects on pulse evolution were studied using pulses with characteristic frequencies that corresponded to spectral wave components ( $k$ ) located in BG and NGV regions
- The BG and NGV regions did not have a noticeable distorting effect on wave evolution because the characteristic wavelengths of waves were relatively large compared to the microstructure found in the felt

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